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## About one class of continual approximate solutions with arbitrary density

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The kinetic Boltzmann equation is one of the central equations in classical mechanics of many-particle systems. For the model of hard spheres it has a form [1, 2]:

$$D(f) = Q(f, f). \quad (1)$$

We will consider the continual distribution [3]:

$$f = \int_{\mathbb{R}^3} du \int_0^{+\infty} d\rho \varphi(t, x, u, \rho) M(v, u, x, \rho), \quad (2)$$

which contains the local Maxwellian of special form describing the screw-shaped stationary equilibrium states of a gas (in short-screws or spirals). They have the form:

$$M(v, u, x, \rho) = \rho \left( \frac{\beta}{\pi} \right)^{\frac{3}{2}} e^{-\beta(v-u-[\omega \times x])^2}. \quad (3)$$

Physically, distribution (3) corresponds to the situation when the gas has an inverse temperature  $\beta = \frac{1}{2T}$  and rotates in whole as a solid body with the angular velocity  $\omega \in R^3$  around its axis on which the point  $x_0 \in R^3$  lies,

$$x_0 = \frac{[\omega \times u]}{\omega^2}, \quad (4)$$

The square of this distance from the axis of rotation is

$$r^2 = \frac{1}{\omega^2} [\omega \times (x - x_0)]^2, \quad (5)$$

$\rho$  is the arbitrary density,  $u \in R^3$  is the arbitrary parameter (linear mass velocity for  $x$ ), for which  $x||\omega$ , and  $u + [\omega \times x]$  is the mass velocity in the arbitrary point  $x$ . The distribution (3) gives not only a rotation, but also a translational movement along the axis with the linear velocity

$$\frac{(\omega, u)}{\omega^2} \omega,$$

Thus, it really describes a spiral movement of the gas in general, moreover, this distribution is stationary (independent of  $t$ ), but inhomogeneous.

The purpose is to find such a form of the function  $\varphi(t, x, u, \rho)$  and such a behavior of all hydrodynamical parameters so that the uniform-integral remainder [3]

$$\Delta = \sup_{(t,x) \in \mathbb{R}^4} \int_{\mathbb{R}^3} |D(f) - Q(f, f)| dv, \quad (6)$$

and its modification "with a weight":

$$\tilde{\Delta} = \sup_{(t,x) \in \mathbb{R}^4} \frac{1}{1 + |t|} \int_{\mathbb{R}^3} |D(f) - Q(f, f)| dv, \quad (7)$$

become vanishingly small.

Also some sufficient conditions to minimization of remainder  $\Delta$  and  $\tilde{\Delta}$  are found. In this work we succeeded a few to generalize results, which obtained in [3]. The obtained results are new and may be used with the study of evolution of screw and whirlwind streams.

1. *Cercignani C.* The Boltzman Equation and its Applications. – New York: Springer, 1988.
2. *Kogan M. N.* The dynamics of a Rarefied Gas. – Moscow: Nauka, 1967.
3. *Gordevskyy V. D., Sazonova E. S.* Continual approximate solution of the Boltzmann equation with arbitrary density // *Matematychni Studii.* – 2016. – Vol. 45, № 2. – P. 194–204.

## **ПРО ОДИН КЛАС КОНТИНУАЛЬНИХ РОЗПОДІЛІВ З ДОВІЛЬНОЮ ГУСТИНОЮ**

*Побудовано новий клас явних наближених розв'язків нелінійного рівняння Больцмана для моделі твердих куль. Він має вид континуальної суперпозиції локальних максвелівських мод, що описують гвинтоподібні стаціонарні рівноважні стани газу з довільною густиною. Отримані деякі граничні випадки, в яких цей розподіл мінімізує інтегральний відхил та відмінний від нього інтегральний відхил з вагою між частинами рівняння.*